



Molecular Crystals and Liquid Crystals Incorporating Nonlinear Optics

Publication details, including instructions for authors and subscription information:

<http://www.tandfonline.com/loi/gmcl17>

Time Variation of Energy in Colliding Solitons in Polyacetylene

Kaoru Iguchi^a

^a Department of chemistry, Waseda University, Tokyo, Japan

Version of record first published: 04 Oct 2006.

To cite this article: Kaoru Iguchi (1990): Time Variation of Energy in Colliding Solitons in Polyacetylene, Molecular Crystals and Liquid Crystals Incorporating Nonlinear Optics, 183:1, 303-306

To link to this article: <http://dx.doi.org/10.1080/15421409008047467>

PLEASE SCROLL DOWN FOR ARTICLE

Full terms and conditions of use: <http://www.tandfonline.com/page/terms-and-conditions>

This article may be used for research, teaching, and private study purposes. Any substantial or systematic reproduction, redistribution, reselling, loan, sub-licensing, systematic supply, or distribution in any form to anyone is expressly forbidden.

The publisher does not give any warranty express or implied or make any representation that the contents will be complete or accurate or up to date. The accuracy of any instructions, formulae, and drug doses should be independently verified with primary sources. The publisher shall not be liable for any loss, actions, claims, proceedings, demand, or costs or damages whatsoever or howsoever caused arising directly or indirectly in connection with or arising out of the use of this material.

TIME VARIATION OF ENERGY IN COLLIDING SOLITONS IN POLYACETYLENE

KAORU IGUCHI

Department of chemistry, Waseda University, Tokyo, Japan

Abstract Time variation of energy and number of solitons in colliding process are examined in polyacetylene. It is strongly dependent on the soliton velocity, but weakly on the soliton width.

THEORY

From the canonical transformation of SSH Hamiltonian we have the electronic part H_S as

$$H_S = \sum_k (B + A k^2) a_k^\dagger a_k - \frac{F}{2N} \sum_{k_1, k_2, k_3} a_{k_1-k_2}^\dagger a_{k_2}^\dagger a_{k_3} a_{k_2-k_1} \quad (1)$$

In the coherent state representation the two-soliton function is given by $\alpha_1(x, t) + \alpha_2(x, t)$ where¹

$$\alpha(x, t) = \sqrt{\xi} \operatorname{sech}[2\xi(x - x_0 - 4c\tau)] \exp i[2cx - 4(c^2 - \xi^2)\tau + \theta], \quad (2)$$

$\tau = At/\hbar$ and the total energy as

$$\begin{aligned} E &= \langle \alpha_1 + \alpha_2 | H_S | \alpha_1 + \alpha_2 \rangle \\ &= \frac{Nd}{2\pi} \int dq n_q (B + A q^2) - \frac{F}{2d} \int dx |\alpha_1(x, t) + \alpha_2(x, t)|^4 \quad (3) \\ &= BN_S + AN_k - \frac{F}{4d\xi} \text{SSC or SAC}, \quad \xi = 1/(2\zeta) \end{aligned}$$

where N_S , N_k and SSC(SAC) correspond to soliton number, kinetic and correlation energies. Two solitons are located at $x = \pm x_0$ at $\tau = 0$, and collide at $\tau_0 (= x_0/4c)$. $c_1 = -c_2 = c$. We consider two kinds of pair; soliton-soliton(SS), and soliton-antisoliton(SA). Correspondingly we put $N_S = \text{SSN or SAN}$, $N_k = \text{SSQ or SAQ}$ in Figures.

RESULT AND DISCUSSION

By numerical computation we examined the constancy of energy and number of solitons in course of time $0 \rightarrow \tau_0$. The behaviors of static (number), kinetic and correlation energies

are shown in Table I and Figures 1-3. It is found that the time variation is strongly dependent on $AL (= 2c\xi \sim \text{soliton velocity} \times \text{width})$, and small for large AL , but remarkable for small AL . Also it is very weakly dependent on $GD (= \xi/d \sim \text{soliton width})$.

The reason that the constancy of energy does not hold will be because $|\alpha_1 + \alpha_2\rangle$ is not an exact eigenfunction of H_s . The employed parameters are not exact enough, which may cause some computational error. Also the separation of electronic and lattice parts is not exact quantum-mechanically, being of the order of square of electron phonon coupling through which the surplus or deficient energy is transferred to and from the lattice.

The large soliton velocity means the small electron-phonon coupling. In this case the soliton is almost an exciton which runs scattering phonons. Contrary, the small soliton velocity means the large electron phonon coupling. In this case the soliton is almost a trapped polariton. Some authors^{2,3} have estimated the soliton velocity in polyacetylene as $v_{\text{sol}} \lesssim 2.6v_{\text{sound}}$ and $v_{\text{sound}} = 1.85 \times 10^6 \text{ cm/s}$. So, we take $v_{\text{sol}} = 2.0v_{\text{sound}}$, then we have $AL = 1.06 (GD = 7)$, $0.75 (GD = 5)$ which lie in the region of our assumption; $AL = 3, 1, 0.5$. By the scheme of Ruckh et al.⁴ they have $v_{\text{sol}} < v_{\text{sound}}$. This is the strong coupling case.

Table I. Variations of total energy E and soliton number N_s with time. E in 2.50 eV. $N_s = 2.0$ at $t = 0$.

AL	GD	E(SS)			N_s	E(SA)			N_s
		$t/\tau_0 = 0$	0.9	1.0		0	0.9	1.0	
3	3	6.010	5.933	2.00	2.00	6.010	5.933	2.00	2.00
	5	5.344	5.273	2.00	2.00	5.344	5.270	2.00	2.00
	7	5.166	5.109	2.00	2.00	5.166	5.103	2.00	2.00
1	3	5.050	5.573	2.54	2.54	5.050	4.370	1.55	1.55
	5	4.999	5.892	2.54	2.54	4.980	3.954	1.55	1.55
	7	4.988	6.045	2.54	2.54	4.992	3.809	1.55	1.55
0.5	3	4.960	6.911	3.37	3.37	4.960	2.821	0.63	0.63
	5	4.964	7.612	3.37	3.37	4.969	2.138	0.63	0.63
	7	4.978	7.897	3.37	3.37	4.968	1.899	0.63	0.63

Fig.1 SOLITON NUMBER

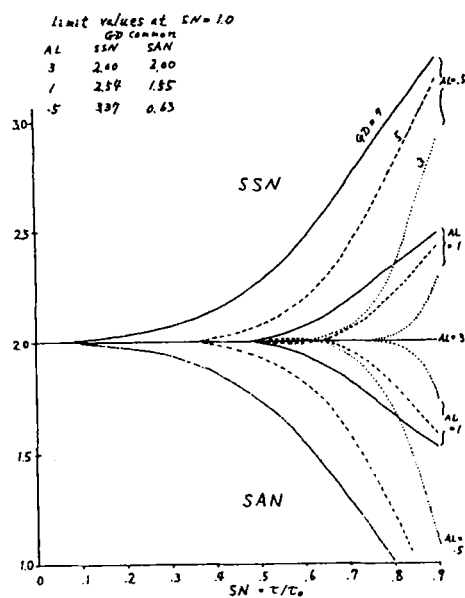
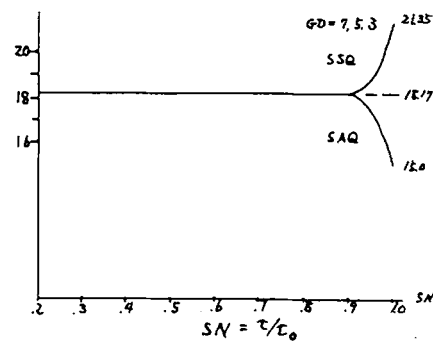
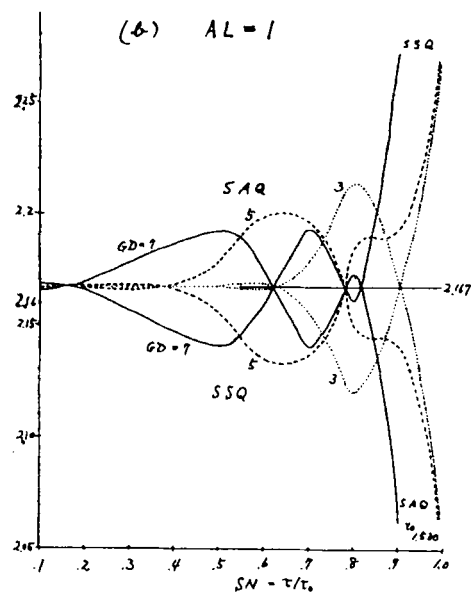


Fig.2 KINETIC ENERGY

(a) $AL = 3$



(b) $AL = 1$



(c) $AL = 0.5$

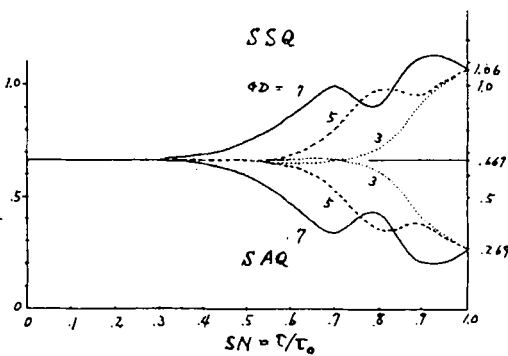
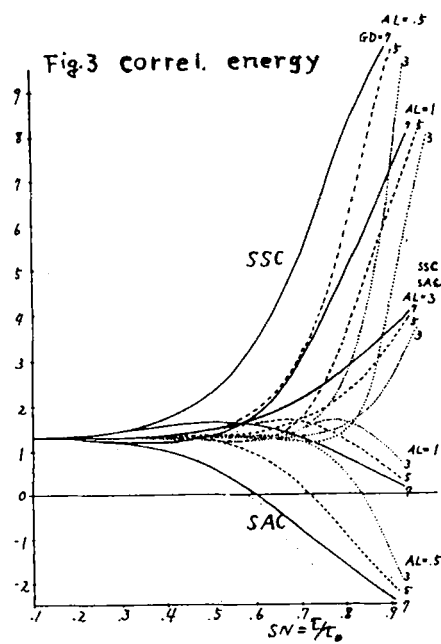


Fig.3 correl. energy



REFERENCES

1. K. Iguchi and H. Kato, Chem. Phys., 88, 315 (1984).
2. F. Guinea, Phys. Rev. B30, 1884 (1984).
3. W. Förner, C.L. Wang, F. Martino and J. Ladik, Phys.Rev. B.37, 4567 (1988).
4. R.Ruckh, E. Sigmund and H. Eisele, J. Chem. Phys. 90, 6463 (1989).